

# Prediction of CO<sub>2</sub> Saturation Spatial Distribution Using Geostatistical Inversion of Time-Lapse Geophysical Data

Dario Grana<sup>1</sup>, Mingliang Liu, and Mohit Ayani

**Abstract**—Carbon dioxide sequestration in deep saline aquifers and depleted reservoirs relies on numerical models for the prediction of the spatial distribution of CO<sub>2</sub> saturation during injection and migration. Due to the limited knowledge of the rock and fluid properties before injection, model predictions are often uncertain and must be updated when new measurements are available. The spatial distribution of CO<sub>2</sub> saturation and the plume location can be monitored using time-lapse geophysical data, such as seismic and controlled source electromagnetic surveys. We propose a geostatistical inversion approach for the prediction of the time-dependent spatial distribution of CO<sub>2</sub> saturation from geophysical data. The methodology is based on the application of a stochastic optimization method, the Ensemble Smoother, for the solution of the inverse problem, using rock physics and geophysical models. The inversion is applied to the difference in the geophysical data acquired before and during injection. The predicted models of CO<sub>2</sub> saturation are obtained by updating an ensemble of geostatistically generated prior realizations, based on the misfit between geophysical model predictions and measured data. The novelty of the approach is the integration of geostatistical algorithms and stochastic optimization methods for the joint inversion of geophysical data. The proposed approach allows including hydrological constraints in the prior model and quantifying the prediction uncertainty due to the noise and resolution of the data and approximations in the physical relations. The method is applied to the Johansen formation model, offshore Norway, using synthetic seismic and electromagnetic data.

**Index Terms**—CO<sub>2</sub> sequestration, inverse problems, reservoir geophysics, rock physics, stochastic methods.

## I. INTRODUCTION

**M**ONITORING of injection and migration of CO<sub>2</sub> in deep saline aquifers requires accurate and precise predictions of the temporal–spatial distribution of CO<sub>2</sub> and water, that is, the saturation values of CO<sub>2</sub> and water at each location in the reservoir and their changes through time [1]. If the distributions of CO<sub>2</sub> and water are known at a given time during injection or migration, their distribution at any

future time can be predicted by simulating the fluid flow in the aquifer, according to the injection parameters and the petrophysical properties of the porous rocks, such as porosity and permeability. Fluid flow models and stochastic simulations are commonly used in geoscience applications to predict fluid displacement in porous rocks, as in hydrogeology [2] and energy resources [3]. However, the rock and fluid properties in the subsurface are generally unknown due to the lack of direct measurements that are generally available at the borehole locations only. Therefore, the predictions of fluid saturations are generally uncertain and often inaccurate. The rock and fluid properties can be predicted and updated using the available geophysical data (seismic, electromagnetic, or gravity data) by solving inverse problems based on geophysical models, such as rock physics, seismic wave propagation, and electromagnetic equations [4].

Carbon capture, utilization, and storage in deep saline aquifers have been widely studied [1], [5], and [6]. Several studies focus on geological and geophysical methods to quantify the capacity of the storage unit and fluid flow simulation to predict the CO<sub>2</sub> plume location and the pressure front extent [7]–[12]. Geophysical surveys, in particular reflection seismic data, have often been used for pre-injection reservoir characterization [13]–[20]. Most of these studies use seismic data to predict the petrophysical properties, primarily porosity. Controlled source electromagnetic data are generally more sensitive to fluid volumes than seismic data and are commonly used to map fluid saturations in the subsurface [21]–[30]. The joint inversion of seismic and electromagnetic data has been proposed in the recent geophysical literature [31]–[37]. However, a comprehensive workflow for updating saturation models and for uncertainty quantification based on fluid flow simulations and monitoring geophysical data is still missing.

The estimation of CO<sub>2</sub> and water saturation based on geophysical measurements during injection and migration is an inverse problem where the model variables are the time-dependent saturations in the reservoir and the data are the time-lapse geophysical measurements (i.e., seismic and electromagnetic responses). The goal is to predict the saturation values at any time and location in the reservoir and quantify their uncertainty. Several mathematical algorithms have been presented, including stochastic methods [17], [38], and [39]. In time-dependent problems, such as subsurface model updating with time-dependent data, data assimilation methods are

Manuscript received June 21, 2020; revised July 29, 2020; accepted August 20, 2020. (Corresponding author: Dario Grana.)

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Digital Object Identifier 10.1109/TGRS.2020.3018910

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generally used to update the reservoir model. The ensemble-based methods represent a particular category of algorithms in which an ensemble of models is simultaneously updated to match the observed measurements. The ensemble-based methods include filtering and smoothing algorithms such as Ensemble Kalman filter and the Ensemble Smoother [40]–[43] and have been applied to borehole measurements and geophysical data [44]–[50].

In this work, we propose to apply the ensemble-based methods to fluid saturation predictions from time-lapse geophysical data. The proposed inversion approach combines geostatistical methods for the generation of the prior models and a stochastic optimization algorithm for the updating of the models conditioned on the geophysical data. To generate initial model realizations that preserve the hydrological realism of the saturation distributions, we propose to generate the prior realizations by sampling from a large ensemble of fluid flow simulations with variable petrophysical properties, and we update them based on the data misfit. The method is validated on the Johansen formation, offshore Norway. The reservoir model has been presented in [14] and [15]. The synthetic time-lapse geophysical data set includes seismic and controlled source electromagnetic (CSEM) data. The geostatistical approach to geophysical inversion allows predicting accurately the CO<sub>2</sub> and water spatial distributions, and the use of fluid flow simulations for prior models guarantees realistic hydrological features. We first predict the porosity from seismic data (pre-injection survey) assuming that the aquifer is initially saturated with water and then estimate the CO<sub>2</sub> saturation from time-lapse measurements. This study proves the value of the proposed probabilistic modeling method for monitoring the CO<sub>2</sub> plume displacement using geophysical data.

## II. METHODOLOGY

We present the inversion method for the prediction of rock and fluid properties given a set of geophysical measurements.

### A. Problem Setting

In the proposed mathematical notation, the model vector  $\mathbf{m}$  represents the rock and fluid properties (porosity and CO<sub>2</sub> saturation) at each location in the spatial model at a given time and the data vector  $\mathbf{d}$  represents the geophysical measurements, including seismic amplitudes and travel times and electromagnetic amplitudes and phases. The mathematical-physical operator  $\mathbf{f}$  that links the model variables to the data predictions is given by

$$\mathbf{d} = \mathbf{f}(\mathbf{m}) + \mathbf{e} \quad (1)$$

where  $\mathbf{e}$  is the error vector [51]. The operator  $\mathbf{f}$  might include wave propagation equations, Maxwell's equations for the electrical and magnetic fields, and rock physics models to relate porosity and fluid saturations to electrical resistivity and elastic properties. The operator  $\mathbf{f}$  might assume different formulations in different lithologies and rock formations depending on the mineral composition and structure of the porous rocks [4].

The solution of the inverse problem is an approximate model  $\hat{\mathbf{m}}$  that minimizes the mismatch between data and model predictions

$$\hat{\mathbf{m}} = \operatorname{argmin}_m \|\mathbf{d} - \mathbf{f}(\mathbf{m})\|. \quad (2)$$

We adopt a Bayesian approach and compute the posterior distribution  $p(\mathbf{m}|\mathbf{d})$  of the model parameters given the data using Bayes' rule

$$p(\mathbf{m}|\mathbf{d}) = kp(\mathbf{d}|\mathbf{m})p(\mathbf{m}) \quad (3)$$

where  $p(\mathbf{m})$  is the prior distribution of the model parameters,  $p(\mathbf{d}|\mathbf{m})$  is the likelihood associated with the operator  $\mathbf{f}$ , and  $k = 1/p(\mathbf{d}) = 1/\int p(\mathbf{d}|\mathbf{m})p(\mathbf{m})d\mathbf{m}$  is a normalizing constant. In the proposed approach, the error vector  $\mathbf{e}$  is assumed to be distributed according to a multivariate Gaussian distribution  $N(\mathbf{e}; \mathbf{0}, \Sigma_e)$  with  $\mathbf{0}$  mean and known covariance matrix  $\Sigma_e$  and is assumed independent of the model variables. The posterior mean of the distribution  $p(\mathbf{m}|\mathbf{d})$  is the most probable model  $\hat{\mathbf{m}}$  of porosity and CO<sub>2</sub> saturation.

### B. Geophysical Forward Model

We first present the geophysical models used to predict the geophysical response of the aquifer model. The geophysical relations used to predict the seismic and electromagnetic response for all the possible combinations of rock and fluid properties in the model include the following: 1) elastic rock physics model; 2) electrical rock physics model; 3) seismic model; and 4) electromagnetic model.

We assume one mineral phase (e.g, quartz) and two fluid components, water and CO<sub>2</sub>. Therefore, the variables of interest are porosity  $\phi$ , water saturation  $S_w$ , and CO<sub>2</sub> saturation  $1 - S_w$  (where the scalar notation is used to indicate the value at a given location in the model). There are several rock physics models to compute the elastic response (i.e., P-wave velocity, S-wave velocity, and density) of fluid saturated porous rocks with given porosity  $\phi$  and water saturation  $S_w$  [4]. In the proposed approach, we use the soft sand model combined with Gassmann's equations and the density equation. The density of the fluid saturated rock  $\rho(\phi, S_w)$  is computed as a multilinear function of porosity and water saturation

$$\begin{aligned} \rho(\phi, S_w) &= (1 - \phi)\rho_m + \phi\rho_f = \\ &= (1 - \phi)\rho_m + \phi[(1 - S_w)\rho_{\text{CO}_2} + S_w\rho_w] \end{aligned} \quad (4)$$

where  $\rho_m$  is the density of the mineral phase, and  $\rho_f$  is the density of the fluid mixture and depends on water saturation  $S_w$  and the densities of CO<sub>2</sub> and water,  $\rho_{\text{CO}_2}$  and  $\rho_w$ . The P- and S-wave velocities are computed from the elastic moduli and density of the saturated rock. Because we assume one mineral phase, the bulk and shear moduli of the solid are the characteristic values of quartz [4]. In a multimineral setting, approximations of the elastic moduli can be computed using Voigt–Reuss–Hill or Hashin–Shtrikman elastic bounds [4]. The bulk and shear moduli of the dry rock,  $K_{\phi_c}$  and  $G_{\phi_c}$ , at the critical porosity  $\phi_c$  (i.e., the maximum porosity of a

porous rock, beyond which the system can be considered a suspension) are computed using Hertz–Mindlin equations

$$K_{\phi_c} = \sqrt[3]{\frac{c^2(1-\phi_c)^2 G_0^2}{18\pi^2(1-\nu_0)^2} P}$$

$$G_{\phi_c} = \frac{5-4\nu_0}{5(2-\nu_0)} \sqrt[3]{\frac{3c^2(1-\phi_c)^2 G_0^2}{2\pi^2(1-\nu_0)^2} P} \quad (5)$$

where  $c$  is the average number of contacts per grain,  $G_0$  is the shear modulus of the solid phase (at 0 porosity),  $\nu_0$  is the Poisson's ratio of the solid phase, and  $P$  is the effective pressure. The bulk and shear moduli of the dry rock,  $K_d(\phi) = k_{HS^-}(\phi, K_0, K_{\phi_c})$  and  $G_d(\phi) = g_{HS^-}(\phi, G_0, G_{\phi_c})$ , for porosity  $\phi \in [0, \phi_c]$ , are computed using an harmonic average of the elastic moduli of the solid phase ( $K_0, G_0$ ) and the elastic moduli at the critical porosity ( $K_{\phi_c}, G_{\phi_c}$ ), using the functions  $k_{HS^-}$  and  $g_{HS^-}$  representing the modified Hashin–Shtrikman lower bounds [4]. The bulk modulus of the saturated rock,  $K_s(\phi, S_w)$  is computed using Gassmann's equation

$$K_s(\phi, S_w) = K_d(\phi) + \frac{\left(1 - \frac{K_d(\phi)}{K_0}\right)^2}{\frac{\phi}{K_f(S_w)} + \frac{(1-\phi)}{K_0} - \frac{K_d(\phi)}{K_0^2}} \quad (6)$$

where the bulk modulus of the fluid  $K_f(S_w)$  is computed using Reuss average as

$$K_f(S_w) = \left(\frac{(1-S_w)}{K_{CO_2}} + \frac{S_w}{K_w}\right)^{-1} \quad (7)$$

with  $K_{CO_2}$  and  $K_w$  being the bulk moduli of CO<sub>2</sub> and water, respectively, whereas the shear modulus of the saturated rock,  $G_s(\phi, S_w)$ , is equal to the shear modulus of the dry rock  $G_d(\phi)$ , according to Gassmann's theory [4]. The P- and S-wave velocities,  $V_P$  and  $V_S$ , are then computed as

$$V_P = \sqrt{\frac{K_s(\phi, S_w) + 4/3G_s(\phi, S_w)}{\rho(\phi, S_w)}}$$

$$V_S = \sqrt{\frac{G_s(\phi, S_w)}{\rho(\phi, S_w)}} \quad (8)$$

where  $\rho(\phi, S_w)$  is given by (4).

The resistivity  $R$  of the saturated rock can be calculated using Archie's law [4]

$$R = \frac{R_w}{\phi^m S_w^n} \quad (9)$$

where  $R_w$  is the resistivity of water,  $m$  is the cementation exponent, and  $n$  is the saturation exponent. These parameters are assumed to be constant. Archie's law is generally assumed to be valid in sandstone. Other rock physics models such as elastic inclusion models for elastic properties and electrical models accounting for clay inner resistivity [4] could also be applied.

The seismic response can be predicted by solving the poro-elastic wave equation. However, approximated models are also available. In this work, we adopt a convolutional model. Given a sequence of rock formations, their seismogram  $s(\theta)$  can be computed, for any reflection angle  $\theta$ , as a function

of the vectors of the P- and S-wave velocities and density, along the vertical profile. For weak elastic contrasts and small reflection angles, the seismic response  $s(\theta)$  can be accurately approximated as a convolution of the seismic wavelet  $w(\theta)$  and the P–P reflection coefficient series  $r(\theta)$  that depend on the P- and S-wave velocities  $V_P$  and  $V_S$ , and density  $\rho$ , according to Shuey's linearized approximation of Zoeppritz equations [52]. At a given travel time  $t$ , the corresponding seismic amplitude is given by

$$s(t, \theta) = w(t, \theta) * r(t, \theta) = \int w(t, \theta) r(t - u, \theta) du \quad (10)$$

where  $*$  represents the convolution operator.

The CSEM electric field  $\mathbf{E}$  and the magnetic field  $\mathbf{H}$  can be computed based on the electrical resistivity obtained from Archie's law. The relations linking electrical conductivity  $\sigma = 1/R$  (i.e., the reciprocal of resistivity) to the curl of the electric field  $\mathbf{E}$  and curl of the magnetic field  $\mathbf{H}$  are described by Maxwell's equations

$$\begin{aligned} \nabla \times \mathbf{E} - i\omega\mu\mathbf{H} &= \mathbf{M}^S \\ \nabla \times \mathbf{H} - \sigma\mathbf{E} &= \mathbf{J}^S \end{aligned} \quad (11)$$

where  $\omega$  is the angular frequency,  $\mu$  is the magnetic permeability (i.e. the resistance to the magnetic field), and  $\mathbf{M}^S$  and  $\mathbf{J}^S$  are the electric and magnetic sources, respectively. Assuming an isotropic 2-D conductivity model along the strike direction, Maxwell's equations can be solved using finite element methods in the frequency domain, as shown in [25] and [53]. The electrical conductivity, in principle, is complex and depends on the real electrical conductivity and the dielectric permittivity; however, for low-frequency sources, the conductivity is assumed to be equal to the real component. The 2-D formulation is extended to 3-D applications by applying Maxwell's equations section by section, where each section is processed independently.

In our approach, the reservoir properties, including porosity and fluid saturations, are defined in the irregular stratigraphic grid of the structural model and then interpolated on a regular grid for the calculation of the seismic data and on an adaptive triangular mesh for the calculation of the electromagnetic data. The spatial interpolation is considered part of the forward operator  $\mathbf{f}$  in (1).

### C. Inverse Method

Next, we introduce the inverse method. The inversion is divided in two steps: first, we predict porosity from the base seismic survey and, then, we predict CO<sub>2</sub> saturation from time-lapse seismic and electromagnetic data. The inverse method is based on the Ensemble Smoother [40], in which an ensemble of prior realizations is first generated and then updated using a Bayesian updating step based on the Kalman filter equations. We first discuss the prior model generation and then present the updating approach.

1) *Prior Ensemble Generation*: Prior model realizations of porosity and saturation can be generated using geostatistical simulations, such as Sequential Gaussian Simulations or Probability Field Simulations algorithms; however, saturation



distributions are generally non-stationary due to the effects of gravity. Traditional algorithms might lead to unphysical model realizations that do not preserve the physical order of fluids or do not obey to hydrological constraints. In the proposed approach, we generate prior realizations according to two different approaches for porosity and saturations, respectively. For the prediction of porosity from the base seismic survey, we assume that the prior model of porosity is a truncated Gaussian random field with locally variable mean and known spatial covariance matrix. We then generate  $N_e$  prior realizations using the fast Fourier transform moving average (FFT-MA) method [54]. The FFT-MA generates a spatially correlated realization by computing the inverse Fourier transform of the product of the Fourier transform of a spatial filter (associated with the spatial covariance function) and the Fourier transform of a spatially uncorrelated realization of a standard Gaussian random field [54]. The FFT-MA method is extremely efficient to generate unconditional realizations; however, any other geostatistical algorithm could be used. For the prediction of CO<sub>2</sub> saturation from time-lapse data, we first generate  $N_e$  prior realizations of porosity and permeability, using the FFT-MA method, and then, for each realization, we run a fluid flow simulation and choose a saturation model at a random time according to a uniform distribution in the simulation time interval. It is possible to choose the saturation realizations at the time at which the data are measured, with the differences in the realizations being due to the initial porosity and permeability values; however, the variability of such ensemble is limited and the inversion might lead to an underestimation of the posterior uncertainty. Instead, the proposed approach generates a prior ensemble of saturation models with a large variability. The fluid flow simulation is conducted using MATLAB reservoir simulation toolbox (MRST) [55], specifically the MRST-co2lab tool, to mimic CO<sub>2</sub> and water displacement during injection and migration. We store the prior model realizations in the ensemble of vectors  $\{\mathbf{m}_j^{\text{prior}}\}_{j=1,\dots,N_e}$ .

Because volumetric fractions are bounded variables in the interval  $[0,1]$ , we first apply a logit transformation to the model parameters to map porosity and saturations to the (unbounded) set of real numbers, then apply the inversion in the transformed space, and apply the inverse transformation to obtain the final predictions in the porosity and saturation bounded domain.

2) *Ensemble Updating*: The prior realizations are then updated using the Ensemble Smoother approach, where the covariance matrices are approximated using the sample covariance matrices estimated from the ensemble. The initial ensemble includes  $N_e$  models  $\mathbf{m}_j$  for  $j = 1, \dots, N_e$ . The inversion approach can be summarized as follows.

- 1) For each model in the ensemble, we apply the geophysical functions in (4)–(11) to compute the corresponding seismic and electromagnetic response and we obtain the vector of the predicted data  $\{\mathbf{d}_j^{\text{prior}}\}_{j=1,\dots,N_e}$

$$\mathbf{d}_j^{\text{prior}} = \mathbf{f}(\mathbf{m}_j^{\text{prior}}) + \mathbf{e}_j \quad (12)$$

where  $\{\mathbf{e}_j\}_{j=1,\dots,N_e}$  represent the data errors.

- 2) For each model in the ensemble, we compute a stochastic perturbation  $\{\mathbf{d}_{pj}\}_{j=1,\dots,N_e}$  of the measured data  $\mathbf{d}$

as

$$\mathbf{d}_{pj} = \mathbf{d} + \Sigma_e^{1/2} \mathbf{z}_j \quad (13)$$

where  $\mathbf{z}_j \sim N(\mathbf{0}, \mathbf{I}_n)$  is a vector sampled from a  $n$ -variate Gaussian distribution with  $\mathbf{0}$  mean and covariance matrix equal to the  $n \times n$  identity matrix  $\mathbf{I}_n$ , for  $j = 1, \dots, N_e$ , and  $\Sigma_e^{1/2}$  is the square root of the covariance matrix of the measurement errors.

- 3) We then update the ensemble using the Ensemble Smoother updating equation [40]

$$\mathbf{m}_j^{\text{post}} = \mathbf{m}_j^{\text{prior}} + \Sigma_{m,d}^{\text{prior}} (\Sigma_{d,d}^{\text{prior}} + \Sigma_e)^{-1} (\mathbf{d}_{pj} - \mathbf{d}_j^{\text{prior}}) \quad (14)$$

for  $j = 1, \dots, N_e$  and obtain the ensemble of posterior models  $\{\mathbf{m}_j^{\text{post}}\}_{j=1,\dots,N_e}$ . In (14), the matrix  $\Sigma_{m,d}^{\text{prior}}$  is the cross-covariance matrix of  $\mathbf{m}^{\text{prior}}$  and  $\mathbf{d}^{\text{prior}}$  and the matrix  $\Sigma_{d,d}^{\text{prior}}$  is the covariance matrix of the data  $\mathbf{d}^{\text{prior}}$ .

In the ES-MDA [42], steps 1–3 are repeated to assimilate the measured data multiple times and improve the accuracy of the updated predictions, using inflation factors for the covariance matrix of the data errors. In the ensemble-based methods, it is also possible to update the model variables and the data predictions simultaneously; however, in the proposed implementation, we apply the updating to the model variables only [56] and compute the predicted data using the forward model. A sensitivity analysis on the covariance matrix of the error measurement  $\Sigma_e$  is necessary. In our work,  $\Sigma_e$  is assumed to be diagonal, but especially with real time-lapse geophysical data with imperfect repeatability, error correlations should be introduced to model processing errors in the data.

The inversion for porosity and time-dependent saturation can be solved sequentially or jointly. If we assume that before injection, there is only one fluid component (water) in the aquifer, and the base seismic survey is not affected by the fluid; therefore, the porosity model can be predicted from the preinjection seismic survey using the inverse method in (12)–(14), where the model vector  $\mathbf{m}$  includes porosity and the data vector  $\mathbf{d}$  includes the base seismic survey [49]. Then, fluid saturations are predicted from time-lapse geophysical data using the same inverse method, where the model vector  $\mathbf{m}$  includes CO<sub>2</sub> saturation and the data vector  $\mathbf{d}$  includes the difference in base and monitor seismic and electromagnetic surveys.

Because of the large dimension of geophysical data, the Ensemble Smoother approach is often not practical, since it would require a large number of initial realizations to obtain accurate predictions and a reliable quantification of uncertainty. To overcome this problem, we propose to apply a dimensionality reduction method to map the data in a smaller dimensional space and perform the inversion using the Ensemble Smoother in the reduced space [57]. The data reduction can be obtained using traditional methods such as principal component analysis or multidimensional scaling, or using deep learning algorithms, for example, the deep convolutional auto-encoder (DCAE). Deep learning methods



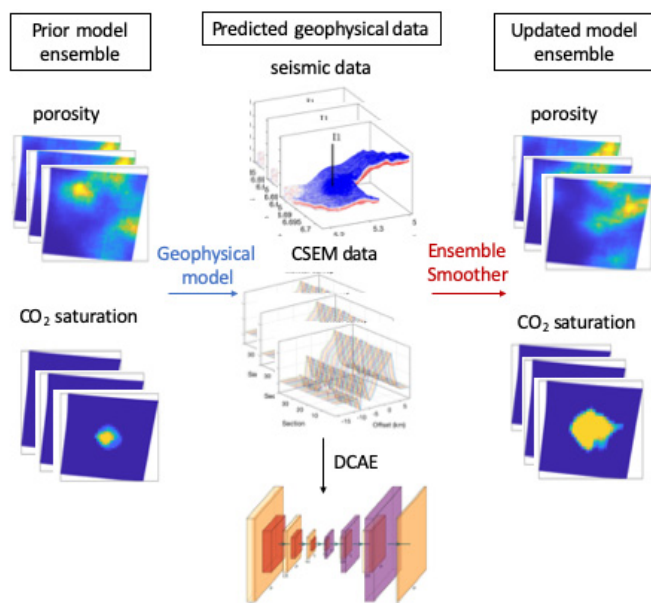


Fig. 1. Schematic of the inversion methodology.

have been successfully applied to seismic inversion [58] and [59]. In the proposed approach, we choose the DCAE approach to compress the data and preserve the spatial correlation structure [57]. The proposed method is schematically depicted in Fig. 1.

### III. APPLICATION

We apply the proposed approach to a carbon dioxide sequestration study based on the Johansen formation model, offshore Norway. The Johansen formation is located below the Troll hydrocarbon field where existing boreholes are available. The data set has been presented in several publications including [14] and [15]. The data show favorable geological features of the storage unit and the overlying sealing, including pore volume, storage capacity, and pressure conditions. The formation depths are between 2200 and 3200 m.

#### A. Synthetic Data Generation

The pre-injection stratigraphic model of the Johansen formation was built based on seismic and well log data. A structural model including porosity and permeability spatial distributions is available [15]. The original geo-cellular model is discretized in  $100 \times 100 \times 5$  cells; however, in this study, we consider a smaller sub-volume of size  $40 \times 40 \times 5$  cells centered around the injection well (Fig. 2). The model includes a major fault interpreted from seismic data. The true reservoir model of porosity and permeability shows relatively high values in the top layers and lower values in the bottom layers. The initial water saturation (before injection) is equal to 1 everywhere. The fluid flow is simulated using MRST-co2lab with an injection period of 100 years and a constant injection rate of  $1.4 \times 10^4$  m<sup>3</sup>/day and migration time of 400 years. Synthetic seismic and CSEM data are computed before injection (base survey) based on the model in Fig. 2 and 10 years after stopping CO<sub>2</sub> injection (monitor survey)

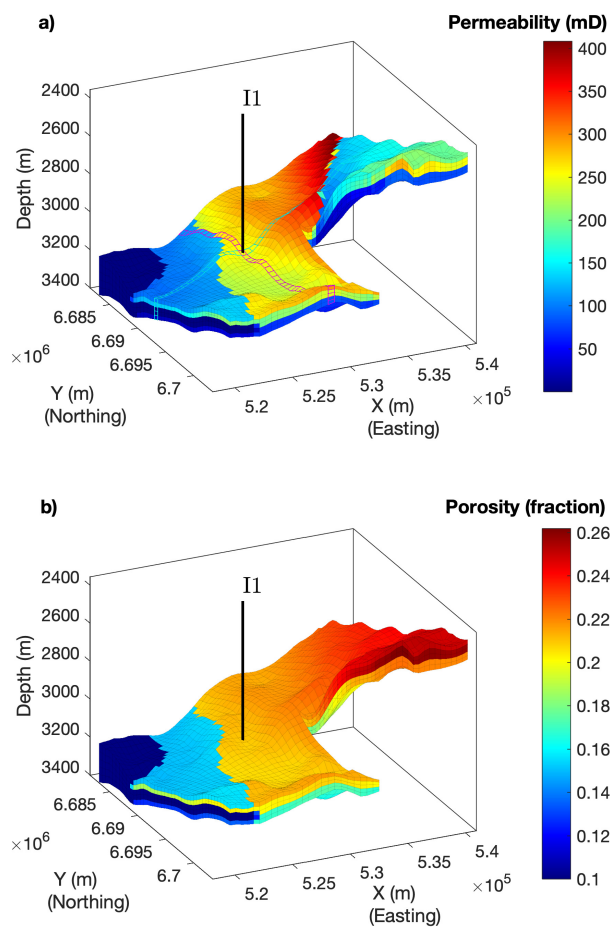


Fig. 2. True reservoir model of (a) permeability and (b) porosity of the deep saline aquifer in the Johansen formation. The coordinates are in the UTM system. The y-axis corresponds to the Northing direction, and the x-axis corresponds to the Easting direction. Black line: Location of the injector well.

based on the initial porosity model and the fluid saturations obtained from the simulator.

The synthetic time-lapse seismic data and CSEM data are generated using the geophysical forward models (4)–(11) and are shown in Figs. 3 and 4. The CSEM amplitudes are computed in the log<sub>10</sub> domain [60] and the differences represent the variations in the logarithm of the amplitudes. Porosity is assumed to be constant in time; therefore, the differences in the geophysical data are due to the changes in saturation. We assume that the effect of pressure changes is negligible compared with the effect of saturation changes. The signal-to-noise ratio of the data is assumed to be 10. The data errors are assumed to be uncorrelated in time and space. Therefore, the covariance matrix of the measurements is diagonal.

#### B. Inversion Results

The inversion methodology is presented in two parts: first, we compute the porosity model based on the base seismic survey; then, we predict the CO<sub>2</sub> saturation based on the time-lapse seismic and CSEM data.

First, we estimate the porosity model from the base seismic survey using the Ensemble Smoother approach. We generate

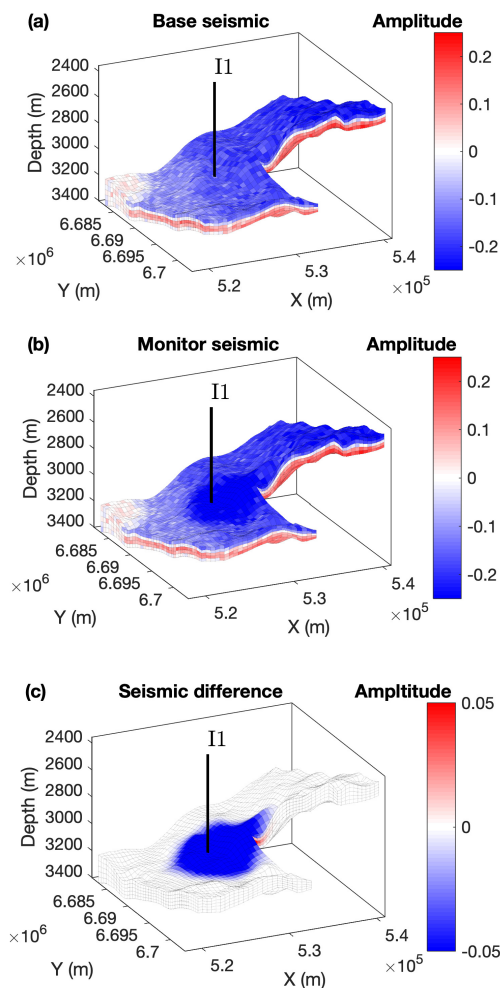


Fig. 3. Time-lapse seismic surveys. (a) Base seismic survey acquired before injection. (b) Monitor seismic survey. (c) Difference between base and monitor seismic data.

the initial ensemble of 100 prior realizations of porosity using the FFT-MA method. The models are generated by adding a locally variable prior mean to preserve the non-stationary trend of porosity in the vertical direction and the spatial anisotropic behavior. The mean of the 100 prior realizations of porosity in the top and bottom layers is shown in Fig. 5. Their seismic responses are generated using the rock physics model and the convolutional seismic model. Due to the large dimension of the data space, we re-parameterize the seismic data using sparse latent features by applying the DCAE and we perform inversion in the lower dimensional space. The posterior mean of the updated porosity models is shown in Fig. 5.

Then, we estimate the CO<sub>2</sub> saturation model in year 110, from monitor seismic and CSEM data. The 100 realizations in the initial ensemble of CO<sub>2</sub> saturation models are generated using dynamic simulations of CO<sub>2</sub> injection and geostatistical simulations of rock properties (permeability and porosity). For each of the geostatistical simulations of porosity and permeability, we simulate the fluid flow after injection using MRST-co2lab and select a saturation realization at a random time according to a uniform distribution on the simulation time interval. Using this approach for the generation of the initial

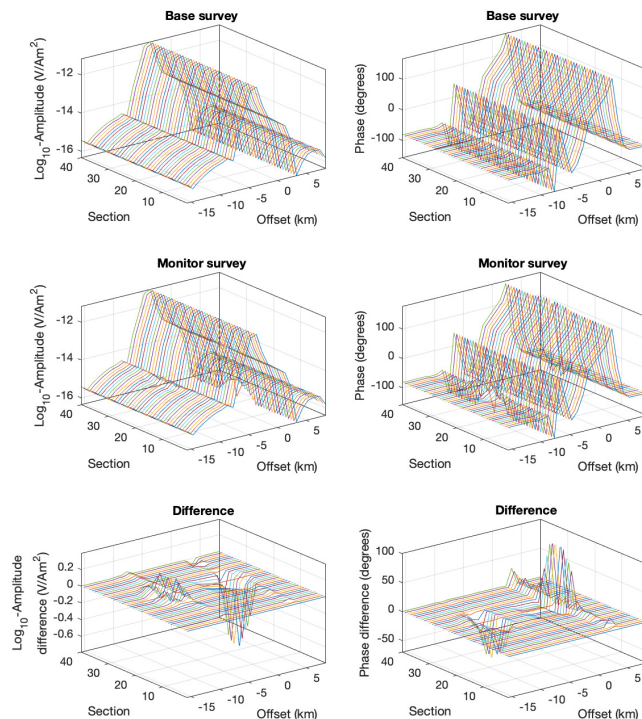


Fig. 4. Time-lapse CSEM surveys acquired before and after injection and their difference. (Left) Log-amplitude data. (Right) Phase. Each 2-D section is represented in a different color for clarity of illustration.

ensemble, we preserve the hydrological realism of the saturation models in the prior and impose physical constraints that cannot be guaranteed by traditional geostatistical simulations. The mean of the CO<sub>2</sub> saturation prior realizations, in the top and bottom layers, is shown in Fig. 6.

The predicted seismic and CSEM responses of the prior realizations in year 110 are obtained by applying the seismic and electromagnetic forward models and using the predicted porosity model obtained in the first step of the inversion. A logit transformation is applied to the saturation models to perform inversion in the real number domain rather than the bounded saturation domain. A dimensionality reduction is applied to data before inversion. We then compute the posterior mean of the CO<sub>2</sub> saturation distribution conditioned on the monitor seismic and CSEM data, using the Ensemble Smoother. The posterior mean of the CO<sub>2</sub> saturation model is shown in Figs. 6 (map view) and 7 (vertical view of two orthogonal sections). The predicted model shows a good agreement with the true model. The standard deviation maps of porosity and CO<sub>2</sub> saturations are shown in Figs. 8 and 9, respectively, which show a reduction in the posterior standard deviation in both properties compared with the prior standard deviation of the initial ensemble.

A 3-D view of the CO<sub>2</sub> plume in year 110 is shown in Fig. 10. The prediction obtained through the geophysical inverse problem accurately matches the true model. To represent the posterior uncertainty in 3-D, we show the 0.90 confidence interval of the CO<sub>2</sub> saturation predictions. Fig. 11 shows the 5<sup>th</sup> and 95<sup>th</sup> percentile of the CO<sub>2</sub> saturation, that is, the lower and upper bounds of the 0.90 confidence interval. Overall, the pointwise confidence intervals are relatively

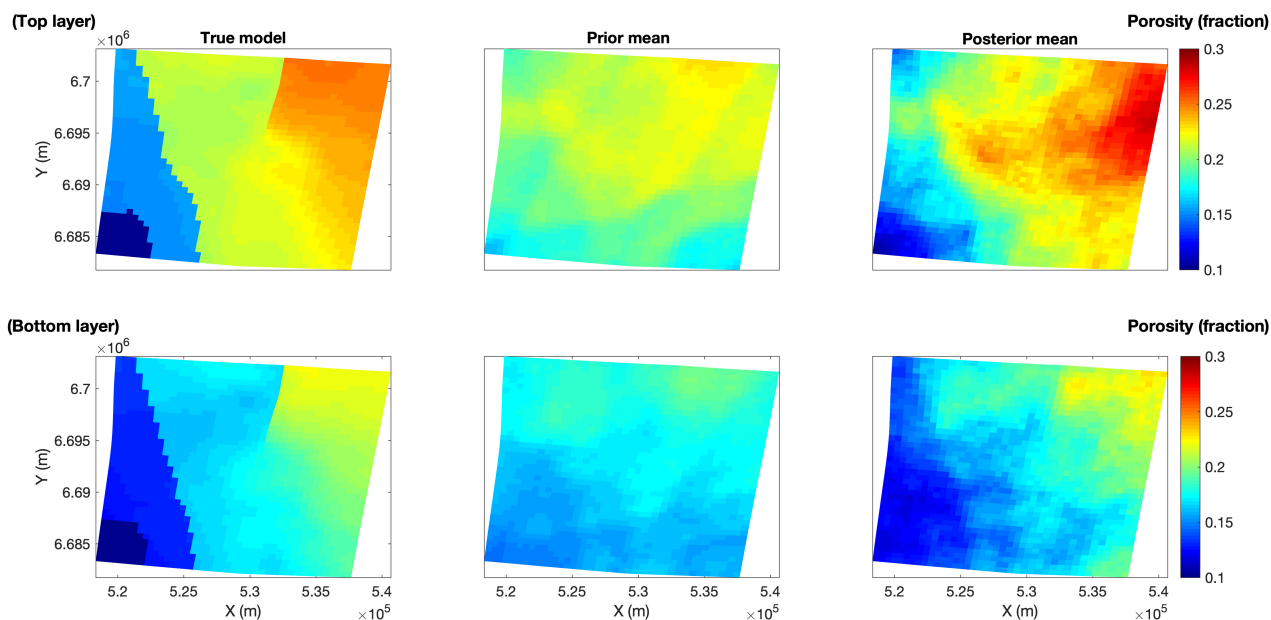


Fig. 5. Porosity model of the top reservoir layer (Top) and the bottom reservoir layer (Bottom). (From Left to Right) True model; prior mean; posterior mean predicted from seismic data.

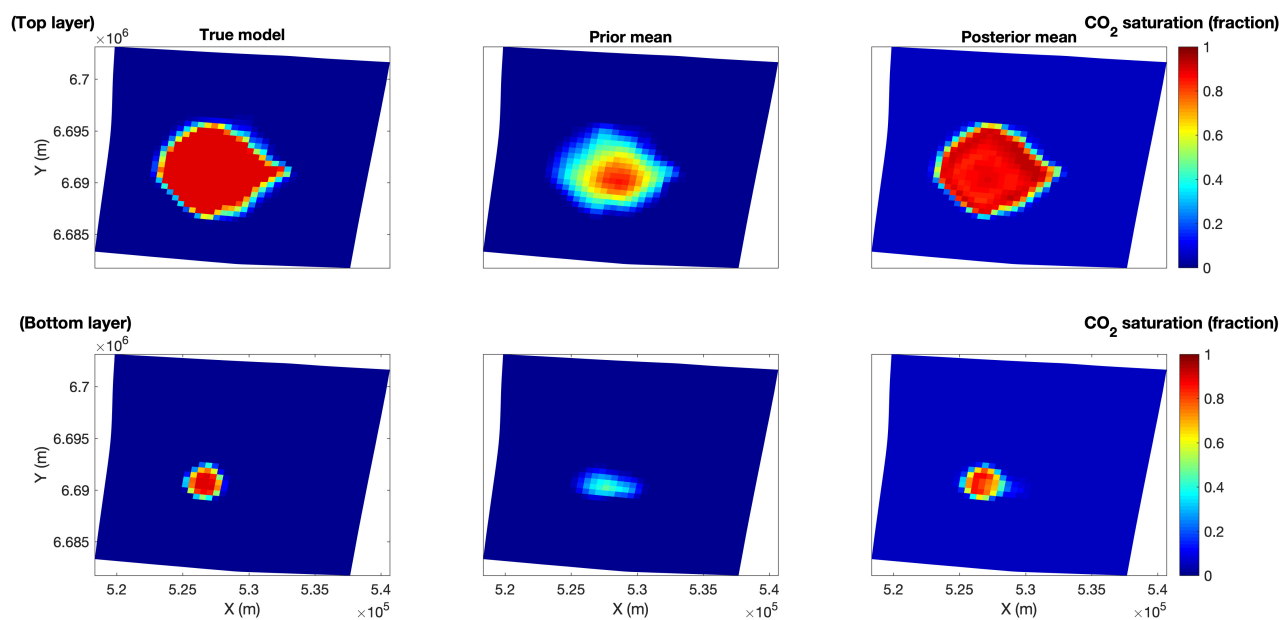


Fig. 6. CO<sub>2</sub> saturation model (map view) in year 110 of the top reservoir layer (Top) and the bottom reservoir layer (Bottom). (From Left to Right) True model; prior mean; posterior mean predicted from time-lapse geophysical data.

narrow, except close to the location near the CO<sub>2</sub> plume where the uncertainty is generally higher. The width of the confidence interval generally depends on the quality of the measured geophysical data and the approximations of the geophysical models. The correlation between the inversion prediction and the true saturation model is 0.96 in the entire model and 0.88 if we only consider the locations where saturation changes during the injection, which shows a high prediction accuracy. To evaluate uncertainty quantification, we computed the percentage of samples of true model within the 0.90 confidence interval, also known as the coverage ratio of the confidence interval, and estimated a value of 83%, which

shows that uncertainty is slightly underestimated, possibly due to model linearization, limited ensemble size, and data reduction.

The proposed inversion results were compared with well-established inverse theory methods such as Occam's inversion and Bayesian linearized inversion which show a higher accuracy and a better reproduction of the spatial correlation structure of CO<sub>2</sub> saturation maps. The studied model is relatively small (8,000 grid cells), and therefore, a limited number of ensemble realizations can be used for stochastic inversion. For a larger model, it might be necessary to increase the number of model realizations in the initial ensemble to avoid the



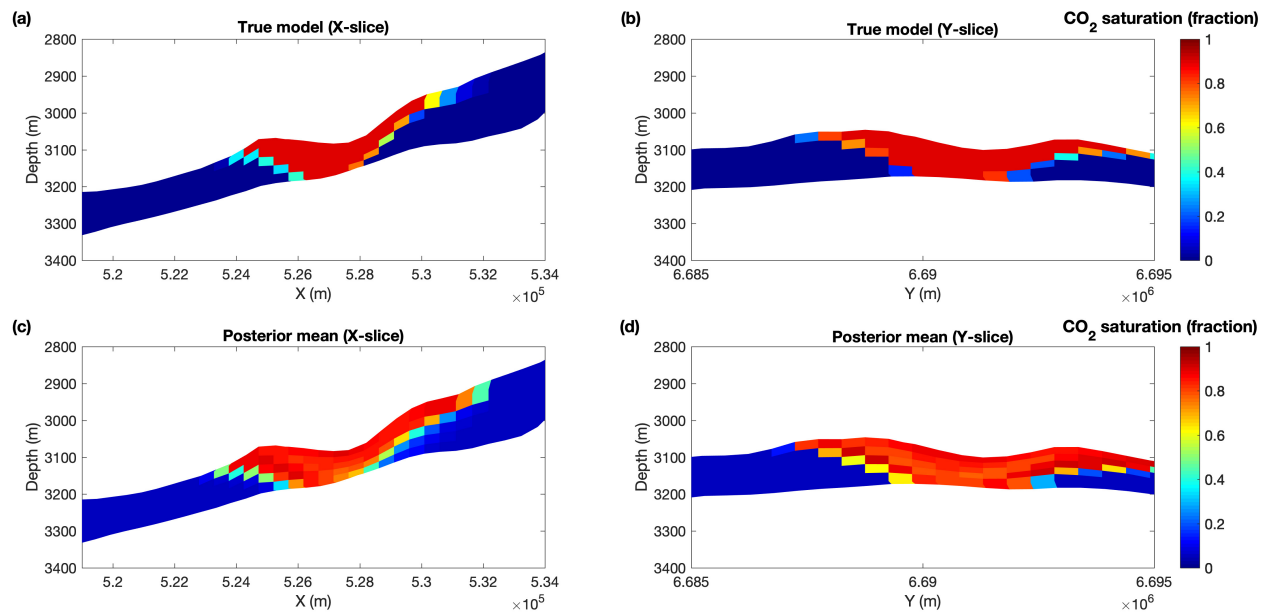


Fig. 7. CO<sub>2</sub> saturation model (section view) in year 110 of two orthogonal sections. (a) True model (X-slice). (b) True model (Y-slice). (c) Posterior mean (X-slice). (d) Posterior mean (Y-slice). The 3-D locations of the X- and Y-slices are shown in Fig. 2(a).

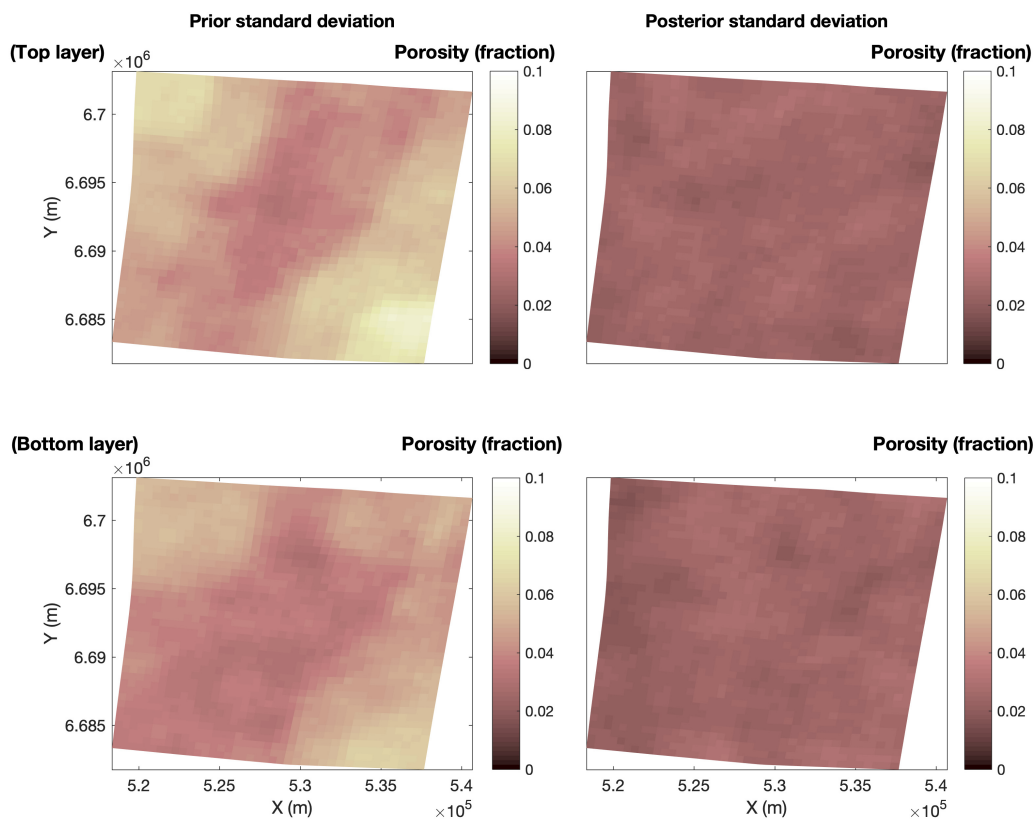


Fig. 8. Uncertainty quantification of the porosity model in the top and bottom layers. (Left) Prior standard deviation maps. (Right) Posterior standard deviation maps.

ensemble collapse. The results generally depend on the quality of the data, in terms of resolution and signal-to-noise ratio. In our application, we assumed a signal-to-noise ratio of 10; however, real data might show lower values and posterior uncertainty might be larger. The dimensionality reduction is generally necessary for 3-D data and the dimension of the

reduced data space can be determined from the magnitude of the eigenvalues of the singular value decomposition of the data vector [49].

The impact of each data set, namely, seismic and CSEM surveys, has been investigated by performing the same inversion using only one of the two data sources. Using only seismic

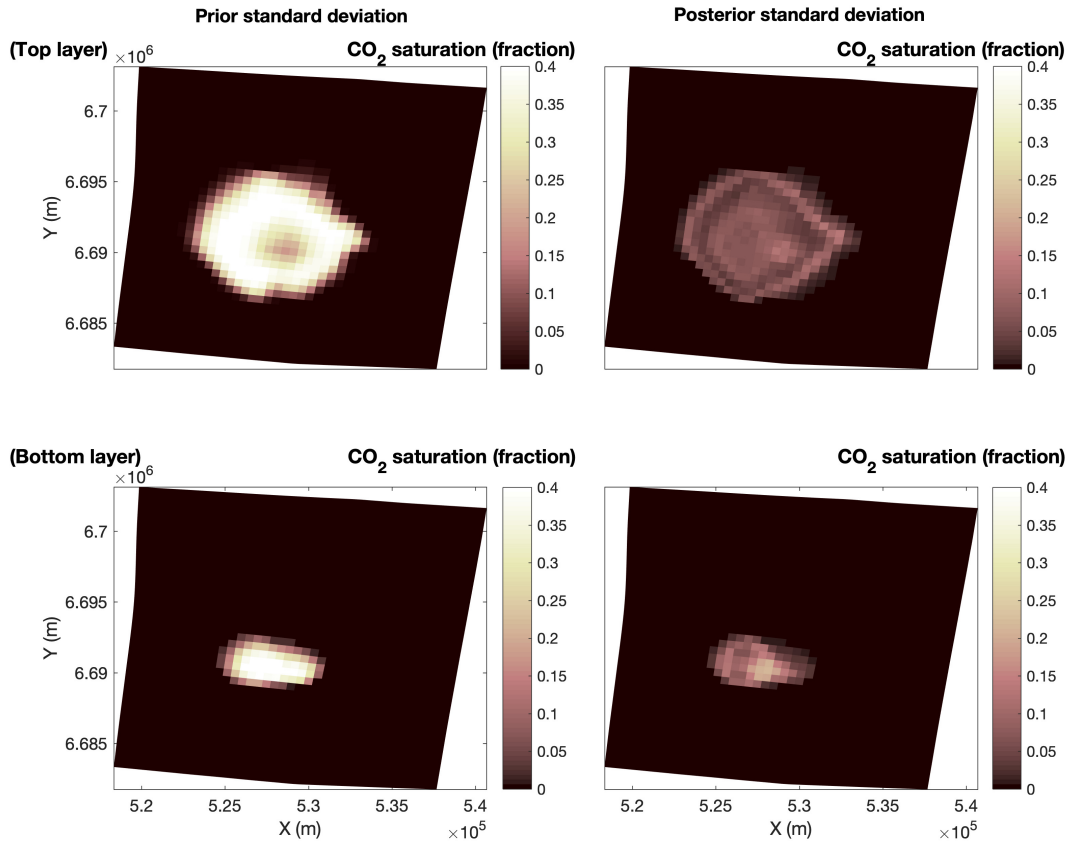


Fig. 9. Uncertainty quantification of the CO<sub>2</sub> saturation model in the top and bottom layers. (Left) Prior standard deviation maps. (Right) Posterior standard deviation maps.

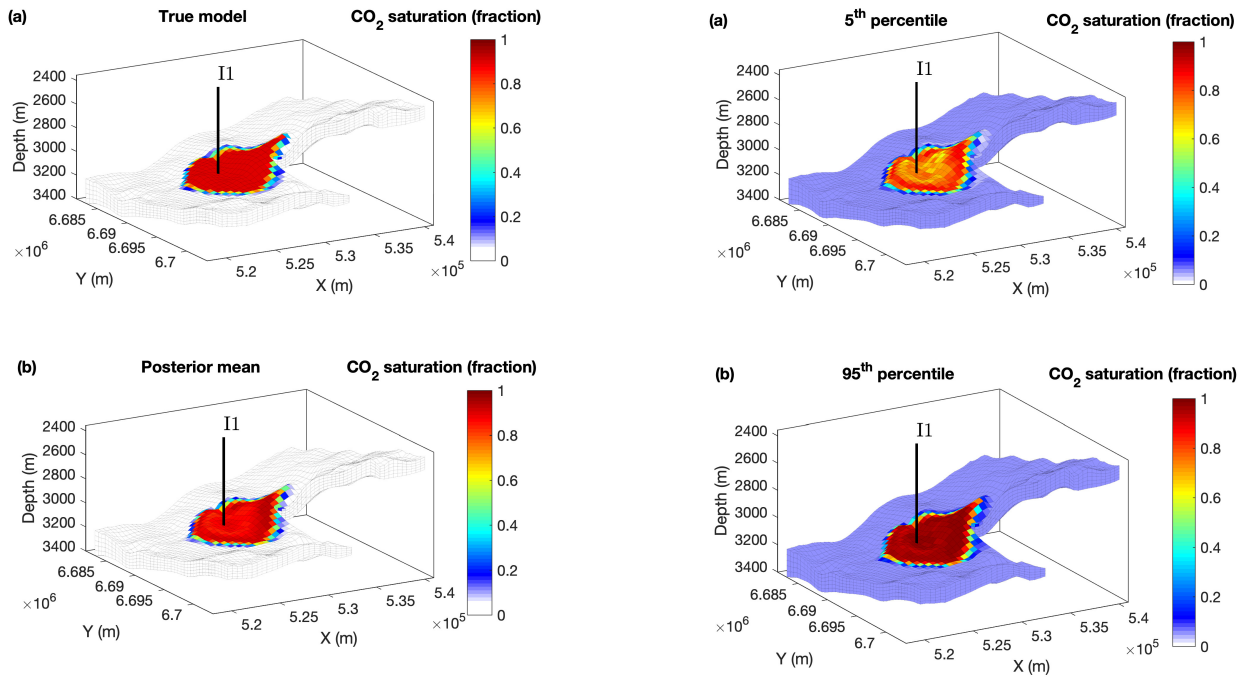


Fig. 10. CO<sub>2</sub> saturation model (3-D view) in year 110. (a) True model. (b) Posterior mean predicted from time-lapse geophysical data.

Fig. 11. Confidence interval 0.90 for the CO<sub>2</sub> saturation model in year 110. (a) 5<sup>th</sup> percentile. (b) 95<sup>th</sup> percentile.

data, inversion can accurately predict the location of the CO<sub>2</sub> plume but the results are less accurate in partially saturated rocks. This is due to the limited sensitivity of velocity to partial saturations and the large uncertainty of density estimation from

seismic data. The correlation between predictions and true saturation model decreases to 0.79. Using only CSEM data, the accuracy of the predictions is strongly dependent on the

frequency of the data, and the posterior standard deviation is overall larger than in the seismic case, due to data resolution. The correlation between predictions and true saturation model decreases to 0.77. Porosity and fluid saturations could be jointly inverted from base and monitor surveys; however, preliminary tests showed that the inversion results overestimate the correlation of the model variables and posterior uncertainty is underestimated.

The inversion code is written in MATLAB, whereas for forward models, we adopted different open source packages including the Fortran code MARE2DEM [53] and the MATLAB code MRST [55]. The computational cost for the inversion is approximately 3 h, with no parallelization, for the proposed model including 16000 model variables (porosity and CO<sub>2</sub> saturation values at each reservoir location) and 14800 measurements (base and monitor seismic and CSEM data) with 74% of the computational time spent on the forward geophysical model.

#### IV. CONCLUSION

We presented a stochastic inversion for the prediction of rock and fluid properties, namely, porosity and time-dependent CO<sub>2</sub> saturation values, from multiple source geophysical data, including seismic and electromagnetic surveys. The proposed Ensemble Smoother approach updates an ensemble of geostatistically generated initial realizations in a Bayesian setting. The method was validated on synthetic time-lapse geophysical data generated for the Johansen formation model. The inversion results show high prediction accuracy and preserve a realistic spatial correlation structure for porosity and CO<sub>2</sub> saturation. This result is achieved by integrating geostatistical simulations to generate the initial porosity models and fluid flow simulations to generate the initial fluid saturation models. The initial realizations are updated conditioned on the mismatch between data predictions and measurements; therefore, the accuracy of the inversion depends on the signal-to-noise ratio and resolution of the data. The posterior standard deviation and confidence intervals provide a quantification of the uncertainty in predictions. The uncertainty in the proposed case study is relatively small due to the use of synthetic data; however, in real data studies, we can expect the saturation uncertainty to increase depending on the quality of the data. The proposed methodology was applied to a deep saline aquifer but could be extended to CO<sub>2</sub> sequestration studies in depleted hydrocarbon reservoirs as well as CO<sub>2</sub> injection and storage as part of EOR applications.

#### ACKNOWLEDGMENT

The authors would like to thank the School of Energy Resources, University of Wyoming and the Nielson Energy Fellowship, for the support. They would also like to thank Dr. K. Key (Lamont-Doherty Earth Observatory, Columbia University) for providing the code MARE2DEM and SINTEF for providing the package MATLAB Reservoir Simulation Toolbox (MRST). The data generated in this project and the inversion codes are available at the following link: <https://tinyurl.com/yaq9y3on>

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